## A stability criterion for coherent states

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joined work with T. Thiemann

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## Motivation

#### Coherent states

- semiclassical limit
- quantum optics
- geometric quantization
- harmonic analysis



#### Why stable?-II

- constrained systems
- group averaging resembles evolution

#### 1 Introduction: Complexifier coherent states

### 2 A stability criterion



## Complexifier coherent states

[Hall '94],[Ashtekar, Lewandowski, Marolf, Mourão, Thiemann '96],[Thiemann, Winkler, Sahlmann '00]

#### Harmonic oscillator

- Complexification:  $(q, p) \mapsto a = q ip$
- Coherent states  $\psi_a$ : eigenstates of  $\hat{a} = \hat{q} i\hat{p}$

 $\mathcal{M}$  Phase space,  $\mathcal{H}$  Hilbert space,  $C \in C^{\infty}(\mathcal{M})$  complexifier,  $\mathcal{L}_C$  Lie derivative along Hamiltonian vector field of C

## Complexifier coherent states

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#### Harmonic oscillator

- Complexification:  $(q, p) \mapsto a = q ip = e^{i\mathcal{L}_C}q$  with  $C = p^2/2$
- Coherent states  $\psi_a$ :  $\psi_a = \left[e^{-\hat{C}}\delta_{q'}(q)\right]_{q' \to a}$

 $\delta_{a'}(q) = \delta(q-q')$  Dirac delta, q' 
ightarrow a analytic continuation

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#### Harmonic oscillator

Complexification: (q, p) → a = q - ip = e<sup>iL<sub>C</sub></sup>q with C = p<sup>2</sup>/2
Coherent states ψ<sub>a</sub>: ψ<sub>a</sub> = [e<sup>-Ĉ</sup>δ<sub>q'</sub>(q)]<sub>q'→a</sub>

 $\delta_{q'}(q) = \delta(q-q')$  Dirac delta, q' 
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#### Complexifier coherent states

A coherent state associated to a complexifier  $C \in C^{\infty}(\mathcal{M})$  is given by  $\psi_z(q) = \left[ e^{-C/\hbar} \delta_{q'}(q) \right]_{q' \to z}$ where  $q' \to z$  denotes analytic continuation to  $z = e^{-i\mathcal{L}_C} q$ .

 $\mathcal{M}$  Phase space,  $\mathcal{H}$  Hilbert space,  $\mathcal{C} \in \mathcal{C}^{\infty}(\mathcal{M})$  complexifier,  $\mathcal{L}_{\mathcal{C}}$  Lie derivative along Hamiltonian vector field of  $\mathcal{C}$ 

## Plan of the talk

#### 1 Introduction: Complexifier coherent states

## 2 A stability criterion

3 Generalized construction

## A stability criterion

#### Definiton

If  $\hat{U}(t, t_0) \psi_{z(t_0)} = e^{i\lambda(t)} \psi_{z(t)} \quad \forall t \in \mathbb{R}^+ \text{ and } \forall z \in \mathcal{M}$ then the system of coherent states  $\{\psi_z\}_{z \in \mathcal{M}}$  is called stable.  $\hat{U}(t, t_0)$  time evolution operator, z(t) classical trajectory in phase space  $\mathcal{M}$ ,  $\lambda(t)$  arbitrary phase

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#### $\Leftrightarrow$ Time-evolution of z(t) depends only on z

$$\Leftrightarrow e^{i\mathcal{L}_{\mathcal{C}}}H = i\,p\,f(q) + g(q)$$

H Hamiltonian, f, g arbitrary functions of q,  $\mathcal{L}_{C}$  Lie derivative along Hamiltonian vector field of C

## First solutions

Want to solve

$$e^{i\mathcal{L}_C}H \stackrel{!}{=} i\,p\,f(q) + g(q)$$

C complexifier, H Hamiltonian, f, g arbitrary functions of  $q, \mathcal{L}_C$  Lie derivative along Hamiltonian vector field of C

## First solutions

Want to solve

$$\sum_{n=0}^{\infty} \frac{i^n}{n!} \{C, H\}_{(n)} \stackrel{!}{=} i \, p \, f(q) + g(q)$$

C complexifier, H Hamiltonian, f, g real functions,  $\{C, H\}_{(n+1)} = \{C, \{C, H\}_{(n)}\}$  multiple Poisson bracket

#### Ansatz

Assume H, C quadratic in momentum and  $\{H, C\}_{(n)} = 0$  for some  $n \in \mathbb{N}$ .

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Radial oscillator  

$$H = \frac{p^2}{2} + \frac{\omega^2}{2}q^2 + \lambda q^{-2} \qquad C = \frac{p^2}{2} + \frac{\lambda}{\omega}q^{-2}$$

$$\lambda, \omega > 0 \text{ some constants}$$

## A No-go theorem

#### After a tedious analysis of the

#### Ansatz

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#### One finds:

#### The radial oscillator and its symplectic transforms are the only solutions.

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### Generalized construction

Dynamics in action-angle coordinates  $(I, \Theta)$ 

$$rac{dI}{dt}=0 \quad ext{ and } rac{d\Theta}{dt}=\omega(I)$$

Suggests to define  
$$y := \sqrt{I}e^{i\Theta}$$

Dynamics: 
$$\frac{d}{dt}y = i \omega(I) y$$
  
Polar decomposition:  $y =: \frac{1}{\sqrt{2}}(Q - iP)$ 

$$\Rightarrow \quad y = \frac{1}{\sqrt{2}} e^{-i\mathcal{L}_C} Q \quad \text{with} \quad C = P^2/2$$

 $\omega(l)$  real function, (Q, P) canonically conjugated,  $\mathcal{L}_{C}$  Lie derivative along Hamiltonian vector field of C

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Wroclaw, September 9

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### Properties

#### Advantages

• complexifier straightforward to construct

#### Disadvantages

- generically complicated dependence of y on (p, q)
- y-parametrization defined only locally

## Free particle

Action-angle coordinates

$$I = \frac{p^2}{2}$$
 and  $\Theta = \frac{q}{p}$ 

#### y-Parametrization

$$y_{\omega} = \frac{p}{\sqrt{2}} e^{i\omega q/p}$$

#### Properties

- locally well-defined except for p = 0
- $\bullet\,$  multi-valued  $\rightsquigarrow$  motion appears to be periodically
- If  $\omega << q/p$  then  $P \approx p + \mathcal{O}(\omega^2)$  and  $Q \approx \omega q + \mathcal{O}(\omega^3)$
- $y_{\omega}$  generators of quasi-periodic functions

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#### Summary

- ✓ stability criterion
- ✓ restricted set of solutions
- ✓ generalized construction

#### Quantum Theory

- which polarization
- global/local structure
- geometric quantization

#### **Stability Criterion**

- approximate stability
- several degrees of freedom

#### Application in LQG

- symmetry reduced models
- 🗞 Loop/standard quantization

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# Thank you!