Tomasz Miller

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Warsaw University of Technology & Copernicus Center (Cracow)





XXXVII Max Born Symposium, Wrocław, 7th July 2016

Tomasz Miller (WUT & CC)

Causality for probability measures

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- Causality (in the context of Lorentzian geometry)
 certain binary relation between events (= spacetime points).
- Pointlike events are **not** directly observable!
 - Measuring apparatus' imperfection
 - Too precise measurement => black hole formation
 - Quantum effects non-locality
- Goal: extend the causality relation to "nonlocal events"
 → (Borel) probability measures.
 - Classical given by densities $d\mu = \frac{\rho}{\int d^n x} d^n x$
 - ullet Quantum the 'modulus-squared' principle $d\mu = |\psi|^2 d^n x$

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The **causal precedence** relation \leq on a given spacetime \mathcal{M} $p \prec q \Leftrightarrow \exists$ future-directed causal curve γ from p to q, or p = q

 $\text{Recall } J^+ := \{(p,q) \in \mathcal{M}^2: \ p \preceq q\}, \quad J^+(p) := \{q \in \mathcal{M}: \ p \preceq q\}$

• Goal: Define $\mu \leq \nu$ for $\mu, \nu \in \mathscr{P}(\mathcal{M})$

• Measures can be spread also in the time-like direction.

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Causality for probability measures

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Let $\mathcal M$ be a spacetime. For $\mu,\nu\in\mathscr P(\mathcal M)$ we define

$$\begin{split} \mu \preceq \nu & \iff & \exists \, \omega \in \mathscr{P}(\mathcal{M}^2) \text{ such that:} \\ \bullet \, \forall_{A \text{ - Borel}} \quad \omega(A \times \mathcal{M}) = \mu(A), \quad \omega(M \times A) = \nu(A), \\ \bullet \, \omega(J^+) = 1. \end{split}$$

Math caveat: J^+ must be a Borel set! Luckily, this is always true.

ullet ω can be called a causal coupling or a causal transference plan.

• For $\mu = \delta_p$, $\nu = \delta_q$, the only coupling is $\omega = \delta_{(p,q)}$ and so $\delta_p \preceq \delta_q$ iff $p \preceq q$.

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Each infinitesimal part of the probability measure should travel along a future-directed causal curve.

 $\omega \text{ couples } \mu \text{ with } \nu$ $\Rightarrow \exists \{\gamma_{p,q}\}_{p \in \text{supp } \mu, q \in \text{supp } \nu}$ $\omega(J^+) = 1$ $\Rightarrow \gamma_{p,q} - \text{causal fut.-dir.}$

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Causality for probability measures

Characterisations of causality



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Characterisations of causality

\mathcal{M} – causally simple (i.e. no causal loops + J^+ closed)

$$\mu \preceq \nu \quad \iff \quad \forall \operatorname{compact} \mathcal{K} \subseteq \operatorname{supp} \mu \quad \mu(\mathcal{K}) \leq \nu(J^+(\mathcal{K}))$$



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Causality for probability measures

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Characterisations of causality

\mathcal{M} – globally hyperbolic (i.e. admits a Cauchy hypersurface)

$$\mu \preceq \nu \quad \Longleftrightarrow \quad \forall \operatorname{Cauchy hypersurface} \, \mathcal{S} \quad \mu(J^+(\mathcal{S})) \leq \nu(J^+(\mathcal{S}))$$



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$$t \mapsto \mu_t \in \mathscr{P}(\mathbb{R}^{n+1}), \qquad d\mu_t = \delta_t \times \|\psi(t,x)\|^2 d^n x.$$

• Is the quantum evolution causal, i.e. $\mu_s \preceq \mu_t$ if s < t?



- Causality of quantum evolution is a Lorentz-invariant notion!
- Dirac equation does yield a causal evolution!
- Evolution driven by $\hat{H} = \sqrt{\hat{p}^2 + m^2}$ is not causal!

• Consequence: localisation vs causal properties of wave packets

- Hegerfeldt's theorem: Evolution driven by $\hat{H} \ge 0$ of a *compactly* supported wave function breaks causality infinite tails appear immediately.
- Our formalism \Rightarrow Acausality is a property of the system and not of a state.
- Beyond the wave packet formalism: QFT, gauge theories, QG?
- C^* -algebra of observables \rightarrow causality in the space of states ("noncommutative spacetimes") [N. Franco, M. Eckstein 2013–2016]

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[R. Penrose, Road to Reality, 2004]

Causality for probability measures

Thank you for your attention!

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