

# Searching for the QCD critical point through power-law fluctuations of the proton density in heavy ion collisions



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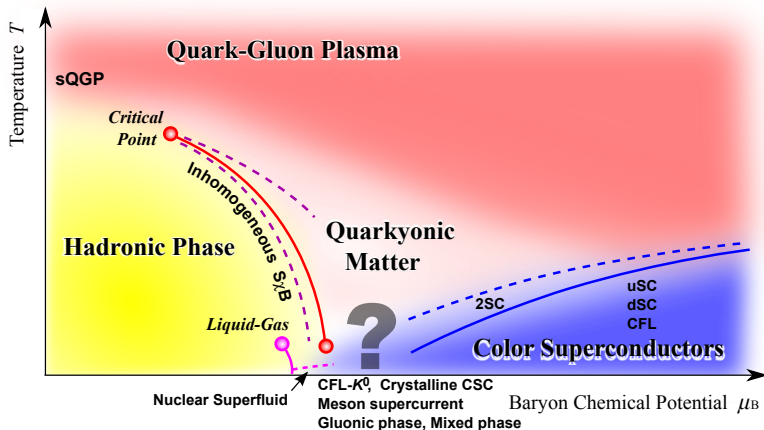
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- 1 QCD Phase Diagram and Critical Phenomena
- 2 Method of analysis
- 3 Results for NA49 data analysis
- 4 NA61 light nuclei feasibility study
- 5 Conclusions and outlook

# Phase diagram of QCD

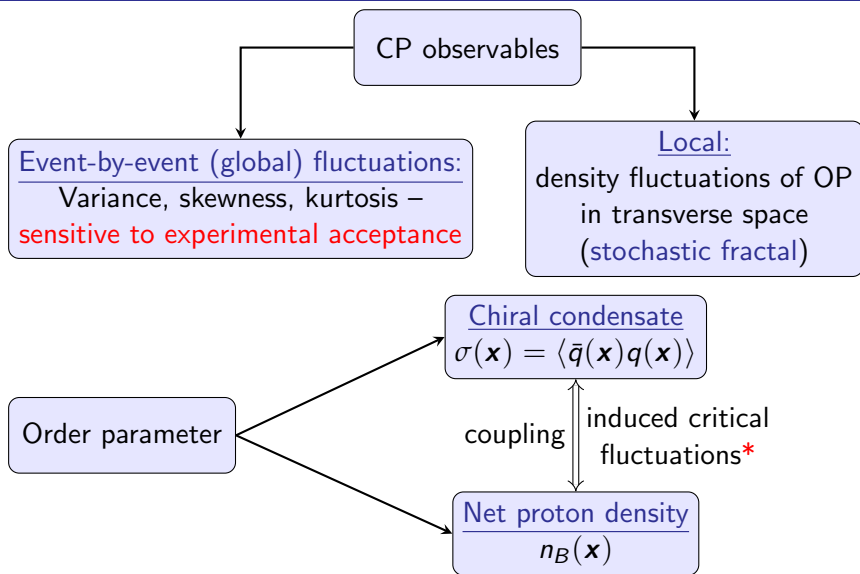
- Objective: Detection / existence of the QCD Critical Point (CP)



*K. Fukushima, T. Hatsuda, Rept. Prog. Phys. 74:014001 (2011)*

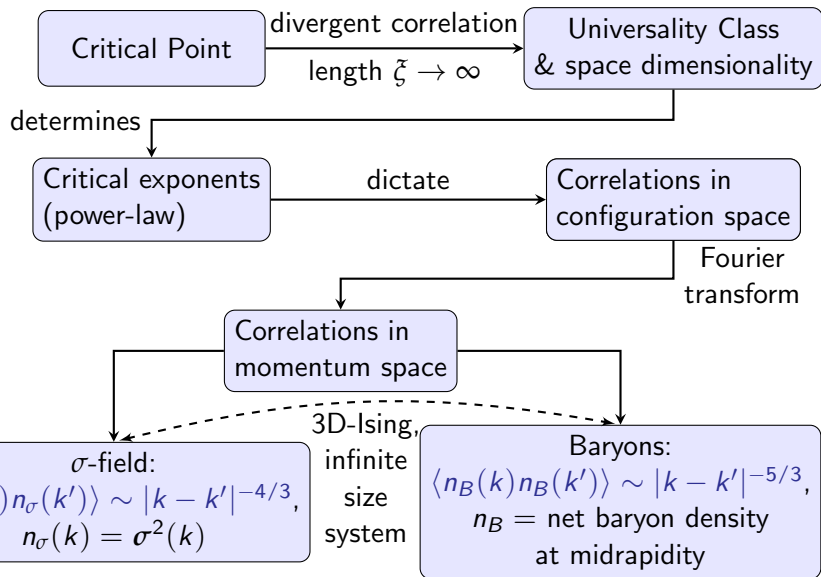
- Look for observables tailored for the CP; Scan phase diagram by varying energy and size of collision system.

# Critical Observables; the Order Parameter (OP)



\*[Y. Hatta and M. A. Stephanov, PRL**91**, 102003 (2003)]

# Self-similar density fluctuations near the CP



# Observing power-law fluctuations

Experimental observation of local, power-law distributed fluctuations



Intermittency in transverse momentum space (net protons at mid-rapidity)

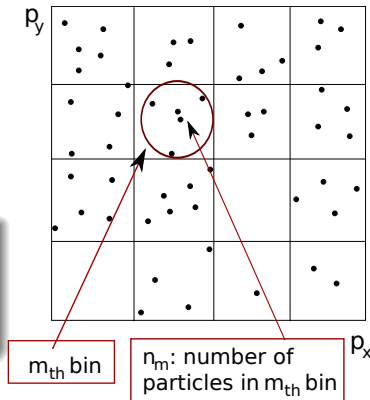
(Critical opalescence in ion collisions\*)

- Transverse momentum space is partitioned into  $M^2$  cells
- Calculate second factorial moments  $F_2(M)$  as a function of cell size  $\Leftrightarrow$  number of cells  $M$ :

$$F_2(M) \equiv \frac{\sum_m \langle n_m(n_m - 1) \rangle}{\sum_m \langle n_m \rangle^2},$$

where  $\langle \dots \rangle$  denotes averaging over events.

\*[F.K. Diakonov, N.G. Antoniou and G. Mavromanolakis, PoS (CPOD2006) 010, Florence]



# Subtracting the background from factorial moments

- Experimental data is **noisy**  $\Rightarrow$  a **background** of uncorrelated/non-critical pairs must be subtracted at the level of factorial moments.
- Intermittency** will be revealed at the level of **subtracted moments**  $\Delta F_2(M)$ .

## Partitioning of pairs into critical/background

$$\langle n(n-1) \rangle = \underbrace{\langle n_c(n_c-1) \rangle}_{\text{critical}} + \underbrace{\langle n_b(n_b-1) \rangle}_{\text{background}} + \underbrace{2\langle n_b n_c \rangle}_{\text{mixed term}}$$

$$\underbrace{\Delta F_2(M)}_{\text{correlator}} = \underbrace{F_2^{(d)}(M)}_{\text{data}} - \lambda(M)^2 \underbrace{F_2^{(b)}(M)}_{\text{background}} - 2 \underbrace{\lambda(M)}_{\text{ratio } \frac{\langle n \rangle_b}{\langle n \rangle_d}} (1 - \lambda(M)) f_{bc}$$

- The **mixed term** can be neglected for dominant background (non-trivial! Justified by **CMC simulations**)

# Scaling of factorial moments – Subtracting mixed events

For  $\lambda \lesssim 1$  (background domination),  $\Delta F_2(M)$  can be approximated by:

$$\Delta F_2^{(e)}(M) = F_2^{\text{data}}(M) - F_2^{\text{mix}}(M)$$

For a critical system,  $\Delta F_2$  scales with cell size (number of cells,  $M$ ) as:

$$\Delta F_2(M) \sim (M^2)^{\varphi_2}$$

where  $\varphi_2$  is the *intermittency index*.

## Theoretical predictions for $\varphi_2$

universality class, effective actions	$\varphi_{2,cr}^{(\sigma)} = \frac{2}{3} \ (0.66\dots)$	$\varphi_{2,cr}^{(p)} = \frac{5}{6} \ (0.833\dots)$
	sigmas (neutral isoscalar dipions)	net baryons (protons)
	[N. G. Antoniou et al, Nucl. Phys. A <b>693</b> , 799 (2001)]	[N. G. Antoniou, F. K. Diakonov, A. S. Kapoyannis, K. S. Kousouris, Phys. Rev. Lett. <b>97</b> , 032002 (2006)]



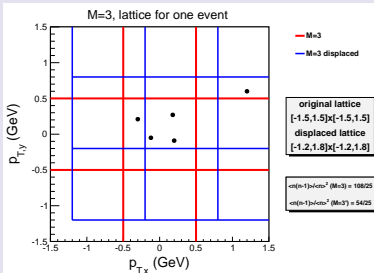
# Statistical & systematic error handling in $F_2(M)$

- $F_2(M)$  averaged over many lattice positions



smoothing of bin boundary effect

## Displaced lattice — a simple example



- Variations of original sample of events produced by resampling (bootstrap) method  $\Rightarrow$  sampling of events with replacement
- $\Delta F_2(M)$  calculated for each bootstrap sample; variance of sample values provides statistical error of  $\Delta F_2(M)$   
[W.J. Metzger, "Estimating the Uncertainties of Factorial Moments", HEN-455 (2004).]
- Distribution of  $\varphi_2$  values,  $P(\varphi_2)$ , and confidence intervals for  $\varphi_2$  obtained by fitting individual bootstrap samples  
[B. Efron, *The Annals of Statistics* 7,1 (1979)]

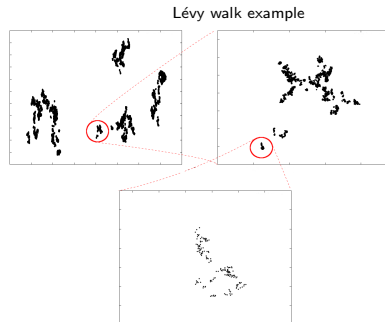
# Critical Monte Carlo (CMC) algorithm for baryons

- Simplified version of CMC\* code:

- Only protons produced
- One cluster per event, produced by random Lévy walk:

$$\tilde{d}_F^{(B,2)} = 1/3 \Rightarrow \phi_2 = 5/6$$

- Lower / upper bounds of Lévy walks  $p_{min,max}$  plugged in.
- Cluster center exponential in  $p_T$ , slope adjusted by  $T_c$  parameter.
- Poissonian proton multiplicity distribution.



## Input parameters

Parameter	$p_{min}$ (MeV)	$p_{max}$ (MeV)	$\lambda_{\text{Poisson}}$	$T_c$ (MeV)
Value	$0.1 \rightarrow 1$	$800 \rightarrow 1200$	$\langle p \rangle_{\text{non-empty}}$	163

\* [Antoniou, Diakonou, Kapoyannis and Kousouris, *Phys. Rev. Lett.* 97, 032002 (2006).]

# NA49 analysed data sets & cuts

- Published in [T. Anticic *et al.*, Eur. Phys. J. C 75:587 (2015), arXiv:1208.5292v5]

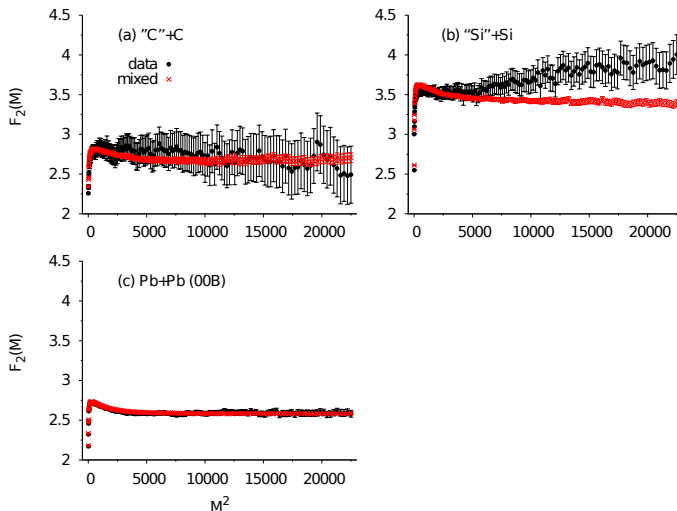
A	"C" + C*	"Si" + Si*	Pb+Pb
# Bootstrap Samples	1000		
Rapidity range	$-0.75 \leq y_{CM} \leq 0.75$		
# lattice positions	11 ( $2 \times 5$ + central)		
Lattice range (GeV)	$[-1.529, 1.471] \rightarrow [-1.471, 1.529]$		
Beam Energy ( $\sqrt{s_{NN}}$ )	158 A GeV (17.3 GeV)		
Centrality range	$0 \rightarrow 12\%$	$0 \rightarrow 10\%$	
Proton purity	$> 80\%$	$> 90\%$	
# events	148 060	165 941	329 789
$\langle p_{data} \rangle$ (after cuts)	$1.6 \pm 0.9$	$3.1 \pm 1.7$	$9.12 \pm 3.15$

\* Beam Components: "C" = C,N, "Si" = Si,Al,P

- Standard NA49 event/track cuts [T. Anticic *et al*, PRC **81**, 149 (2010)].
- $q_{inv}$  cut to remove split tracks, F-D effects and Coulomb repulsion
- Mid-rapidity selected because of approximately constant proton density in rapidity in this region (also avoids nucleons in the corona).  
[N.G. Antoniou, F.K. Diakonos, A.S. Kapoyannis and K.S. Kousouris, PRL **97**, 032002 (2006)]

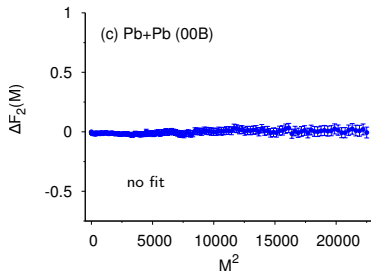
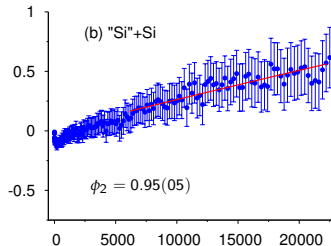
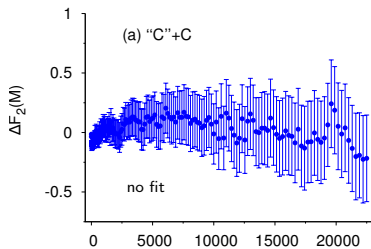
# Analysis results - $F_2(M)$ for protons

- Evidence for intermittent behaviour in “Si” + Si – but large statistical errors.



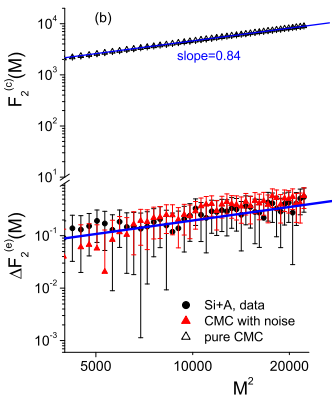
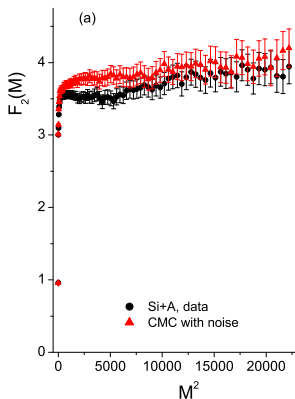
# Analysis results - $\Delta F_2(M)$ for protons

- Fit with  $\Delta F_2^{(e)}(M; \mathcal{C}, \phi_2) = e^{\mathcal{C}} \cdot (M^2)^{\phi_2}$ , for  $M^2 \geq 6000$



# Noisy CMC (baryons) – estimating the level of background

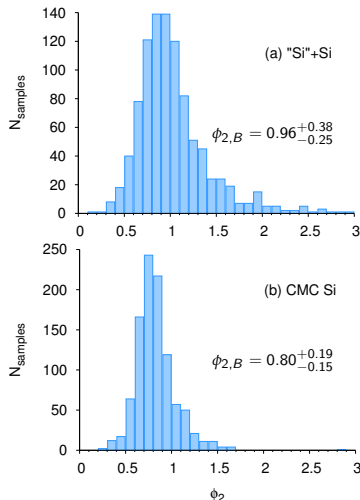
- $F_2(M)$  of noisy CMC approximates “Si” + Si for  $\lambda \approx 0.99$
- Correlator  $\Delta F_2^{(e)}(M)$  has slope  $\phi_2 = 0.80^{+0.19}_{-0.15}$ , very close to  $\phi_2 = 0.84$  of pure  $F_2^{(c)}(M)$



- $\Delta F_2^{(e)}(M)$  reproduces critical behaviour of pure CMC, even though their moments differ by orders of magnitude!
- Noisy CMC results show our approximation is reasonable for dominant background.

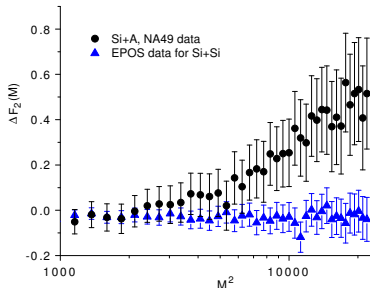
# Analysis results - $\phi_2$ bootstrap distribution

- Distributions are highly asymmetric due to closeness of  $F_2^{(d)}(M)$  to  $F_2^{(m)}(M)$ .



- CMC model with a dominant background can reproduce the spread of  $\phi_2$  values observed in the "Si" + Si dataset
- The spread is partly artificial due to pathological fits (negative  $\Delta F_2(M)$  values in some bootstrap samples)

# Can jets “fake” intermittency effect?



\*[ K. Werner, F. Liu, and T. Pierog,  
Phys. Rev. C 74, 044902 (2006)]

- EPOS event generator\* includes high- $p_T$  jets  $\Rightarrow$  possible spurious intermittency by non-critical protons.
- We simulate 630K Si+Si EPOS events:
  - ①  $Z=14$ ,  $A=28$ , for both beam and target
  - ②  $b_{max} = 2.6$  fm ( 12% most central)
  - ③  $\sqrt{s_{NN}} = 17.3$  GeV
  - ④ Rapidity cuts as in NA49 data
- Intermittency analysis (data & mixed events) repeated for EPOS.
- EPOS clearly cannot account for intermittency presence  $\Rightarrow \Delta F_2(M)$  fluctuates around zero.



# Onwards to NA61 analysis...

- NA49 analysis encourages us to look for intermittency in medium-sized nuclei, in the NA61 experiment.
- Intermittency analysis requires:
  - Large event statistics  $\Rightarrow \sim 100K$  events min., ideally  $\sim 1M$  events.
  - Reliable particle ID  $\Rightarrow$  proton purity should be  $\sim 80\%, 90\%$ .
  - In general,  $5 \rightarrow 10\%$  most central collisions.
  - Adequate mean proton multiplicity in midrapidity ( $\geq 2$ )
- Two candidate NA61 systems – Be+Be @ 150 GeV & Ar+Sc @ 150 GeV.
- What follows is a feasibility study for Be+Be & Ar+Sc – no actual data intermittency analysis at this stage.

# Overview of $^7\text{Be} + ^9\text{Be}$ , $^{40}\text{Ar} + ^{45}\text{Sc}$ @ 150 GeV

## Be+Be:

- Mean proton multiplicity density per event, in mid-rapidity – pilot analysis of NA61 data suggests:

$$\left. \frac{dN_p}{dy} \right|_{|y_{CM}| \leq 0.75, p_T \leq 1.5} \sim 0.75$$

rather low  $\Rightarrow \geq 1.5 \rightarrow 2$  needed

## Ar+Sc:

- Analysis of NA61 data in progress.
- Simulation through EPOS would suggest:

$$\left. \frac{dN_p}{dy} \right|_{|y_{CM}| \leq 0.75, p_T \leq 1.5} \sim 4$$

for

$b_{max} \sim 3.5 \Leftrightarrow \sim 10\%$  centrality;  
adequate for an intermittency analysis

# Simulating Be+Be – EPOS & CMC

- EPOS Simulation parameters:

- ①  ${}^7\text{Be}$  (beam) +  ${}^9\text{Be}$  (target)
- ② Beam energy: 150A GeV (target rest frame)  $\Leftrightarrow \sqrt{s_{NN}} = 16.8$  GeV
- ③ Central collisions  $\Rightarrow b_{\text{max}} = 2.0$  fm
- ④ Total number of simulated events: 200K

## EPOS – proton $p_T$ statistics

$b_{\text{max}}$	#events	$\langle p \rangle_{ p_T  \leq 1.5 \text{ GeV},  y_{CM}  \leq 0.75}$		$\Delta p_{x,y}$
		Non-empty	With empty	
1.0	50,093	$1.48 \pm 0.74$	$0.78 \pm 0.92$	0.43

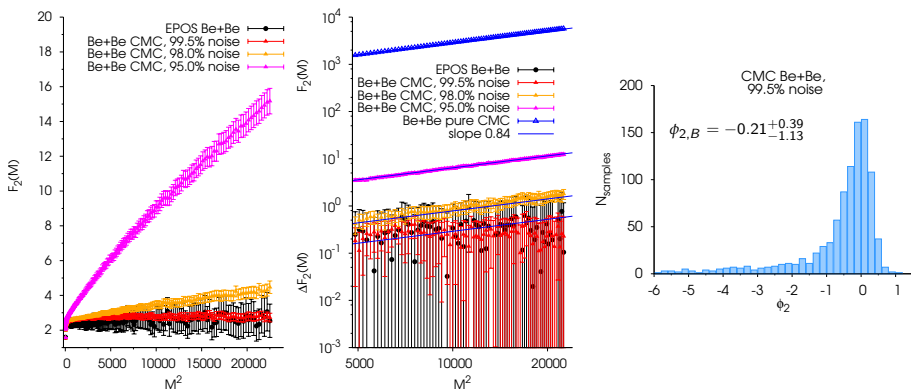
## CMC simulation parameters

Parameter	$p_{\text{min}}$ (MeV)	$p_{\text{max}}$ (MeV)	$\lambda_{\text{Poisson}}$	$T_c$ (MeV)
Value	0.85	1200	0.76	163

- $\langle p \rangle$  in mid-rapidity remains low, except for very central collisions

# Noisy CMC, Be+Be – estimating the level of background

- $F_2(M)$  of noisy CMC approximates Be+Be EPOS for  $\lambda \approx 0.995$
- Correlator  $\Delta F_2^{(e)}(M)$  has slope  $\phi_2 = -0.21_{-1.13}^{+0.39} \rightarrow$  fluctuates around zero.



# Simulating Ar+Sc – EPOS & CMC

- Simulation parameters:

- ①  $^{40}\text{Ar}$  (beam) +  $^{45}\text{Sc}$  (target)
- ② Beam energy: 150A GeV (target rest frame)  $\Leftrightarrow \sqrt{s_{NN}} = 16.8$  GeV
- ③ Central collisions  $\Rightarrow b_{\text{max}} = 3.5$  fm
- ④ Total number of simulated events: 100K

## EPOS – proton $p_T$ statistics

$b_{\text{max}}$	#events	$\langle p \rangle_{ p_T  \leq 1.5 \text{ GeV},  y_{CM}  \leq 0.75}$		$\Delta p_{x,y}$
		Non-empty	With empty	
3.5	100,000	$5.3 \pm 2.5$	$5.3 \pm 2.4$	0.490

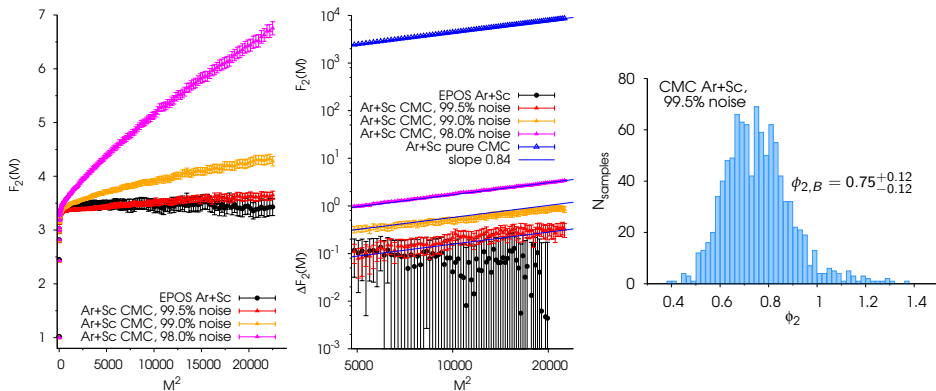
## CMC simulation parameters

Parameter	$p_{\text{min}}$ (MeV)	$p_{\text{max}}$ (MeV)	$\lambda_{\text{Poisson}}$	$T_c$ (MeV)
Value	0.41	1200	5.3	163

- $\langle p \rangle$  in mid-rapidity acceptable for  $b_{\text{max}} \leq 3.5$

# Noisy CMC, Ar+Sc – estimating the level of background

- $F_2(M)$  of noisy CMC approximates Ar+Sc EPOS for  $\lambda \approx 0.995$
- Correlator  $\Delta F_2^{(e)}(M)$  has slope  $\phi_2 = 0.75^{+0.12}_{-0.12}$



Intermittency analysis in transverse momentum space of NA49 data for central “C”+C, “Si”+Si and Pb+Pb collisions has been performed.

- For protons at midrapidity we find significant power-law fluctuations in “Si”+Si at 158A GeV. No significant intermittent behaviour is observed in “C”+C and low-intensity Pb+Pb (00B) data sets.
- The intermittency index  $\phi_2$  for the Si system overlaps with the critical QCD prediction.

# Summary and outlook

- Study of **self-similar (power-law) fluctuations** of the net baryon density provides us with a **promising set of observables** for detecting the location of the QCD critical point.
- **First experimental evidence** for the **approach to the vicinity of the critical point**.
- Analysis favors a CP **close to the freeze-out conditions** of the “Si” + Si system  
*[T. Anticic et al. (NA49 Collaboration), Eur. Phys. J. C 75:587 (2015), arXiv:1208.5292v5]*
- Preliminary study of light nuclei collisions in **NA61 experiment** indicates that an intermittency analysis is **feasible for (at least) the Ar+Sc system** at maximum SPS energy. Performing a systematic intermittency analysis in this system size region (Be+Be, Ar+Sc, Xe+La) will hopefully lead to an **accurate determination of the critical point location**.



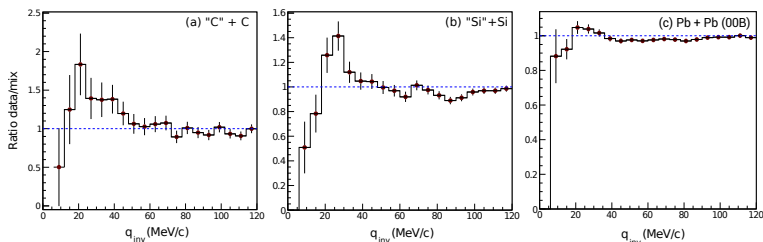
# Thank you!



# Back Up Slides

# Split tracks; the $q_{inv}$ cut in analysed datasets

- Split tracks can create **false positive** for intermittency  $\Rightarrow$  must be **reduced** or **removed**.
- $q_{inv}$ -test – distribution of track pairs:  $q_{inv}(p_i, p_j) \equiv \frac{1}{2} \sqrt{-(p_i - p_j)^2}$ ,  $p_i$ : 4-momentum of  $i^{th}$  track.
- Calculate ratio  $q_{inv}^{data} / q_{inv}^{mixed} \Rightarrow$  **peak** at low  $q_{inv}$  (below 20 MeV/c): **possible split track contamination**.

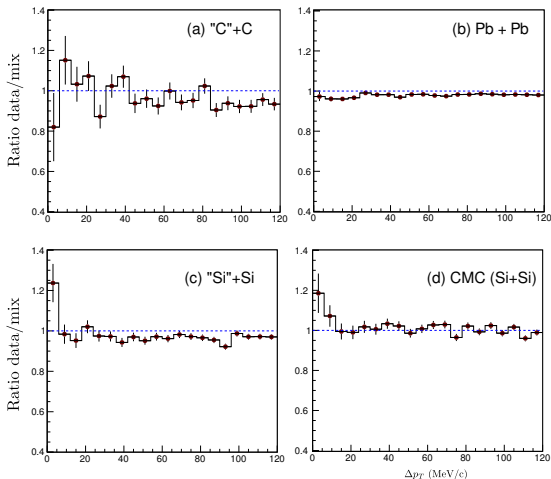


- Anti-correlations due to **F-D effects** and **Coulomb repulsion** must be removed before intermittency analysis  $\Rightarrow$  "dip" in low  $q_{inv}$ , peak predicted around 20 MeV/c [Koonin, PLB 70, 43-47 (1977)]
- Universal cutoff** of  $q_{inv} > 25$  MeV/c applied to all sets before analysis.

# NA49 analysis – $\Delta p_T$ distributions

- We measure correlations in relative  $p_T$  of protons via

$$\Delta p_T = 1/2 \sqrt{(p_{X_1} - p_{X_2})^2 + (p_{Y_1} - p_{Y_2})^2}$$

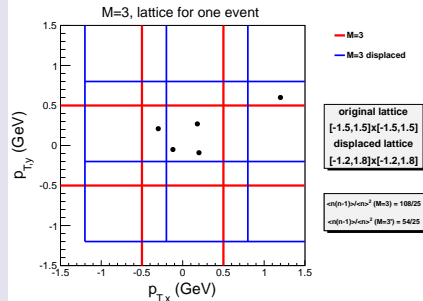


- Strong correlations for  $\Delta p_T \rightarrow 0$  indicate **power-law scaling** of the density-density correlation function  $\Rightarrow$  intermittency presence
- We find a strong peak in the "Si" + Si dataset
- A similar peak is seen in the  $\Delta p_T$  profile of simulated CMC protons with the characteristics of "Si" + Si.

# Improving calculation of $F_2(M)$ via lattice averaging

- **Problem:** With low statistics/multiplicity, lattice boundaries may **split pairs** of neighboring points, affecting  $F_2(M)$  values (see example below).
- **Solution:** Calculate moments several times on **different, slightly displaced lattices** (see example)
- **Average** corresponding  $F_2(M)$  over all lattices. Errors can be estimated by **variance over lattice positions**.
- Lattice displacement is **larger than experimental resolution**, yet **maximum displacement** must be of the order of the **finer binnings**, so as to stay in the correct  $p_T$  range.

## Displaced lattice — a simple example



# Improved confidence intervals for $\phi_2$ via resampling

- In order to estimate the **statistical errors** of  $\Delta F_2(M)$ , we need to produce **variations** of the original event sample. This, we can achieve by using the statistical method of **resampling (bootstrapping)**  $\Rightarrow$ 
  - Sample original events **with replacement**, producing new sets of the **same statistics** (# of events)
  - Calculate  $\Delta F_2(M)$  for each bootstrap sample in the same manner as for the original.
  - The **variance** of sample values provides the statistical error of  $\Delta F_2(M)$ .

[W.J. Metzger, “*Estimating the Uncertainties of Factorial Moments*”, HEN-455 (2004).]

- Furthermore, we can obtain a **distribution**  $P(\phi_2)$  of  $\phi_2$  values. Each bootstrap sample of  $\Delta F_2(M)$  is fit with a power-law:

$$\Delta F_2(M; \mathcal{C}, \phi_2) = e^{\mathcal{C}} \cdot (M^2)^{\phi_2}$$

and we can extract a **confidence interval** for  $\phi_2$  from the distribution of values. [B. Efron, *The Annals of Statistics* 7,1 (1979)]

# Event & track cuts for Si+A

## Event cuts:

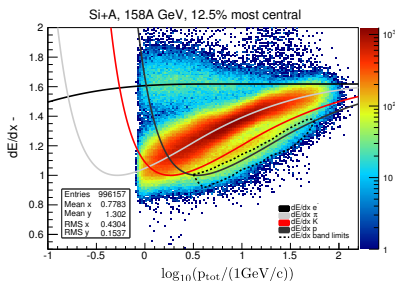
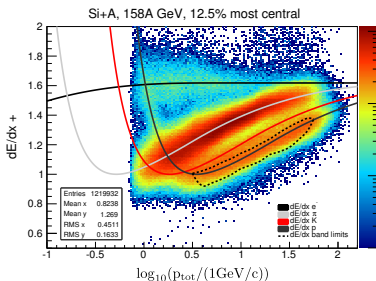
- $\text{lflag} = 0$ ,  $\chi^2 > 0$
- Beam charge cuts (Al,Si,P)
- Vertex cuts:
  - $-0.4 \text{ cm} \leq V_x \leq 0.4 \text{ cm}$
  - $-0.5 \text{ cm} \leq V_y \leq 0.5 \text{ cm}$
  - $-580.3 \text{ cm} \leq V_z \leq -578.7 \text{ cm}$

## Track cuts:

- $\text{lflag} = 0$
- $N_{\text{points}} \geq 30$   
(for the whole detector)
- Ratio  $\frac{N_{\text{points}}}{N_{\text{MaxPoints}}} \geq 0.5$
- $Z_{\text{First}} \leq 200$
- Impact parameters:  
 $|B_x| \leq 2$ ,  $|B_y| \leq 1$
- $dE/dx$  cuts for particle identification
- $p_{\text{tot}}$  cuts (via  $dE/dx$  cut)
- rapidity cut

# NA49 analysis – applied cuts and particle ID

- Cuts based on the standard set of event & track cuts used in NA49 experiment [Anticic et al., PRC,83:054906 (2011)]
- Beam components merged for analysis in “Si” + Si, “C” + C
- Quality cuts to minimize split track effect
- Proton identification through cuts in particle energy loss  $dE/dx$  vs  $p_{TOT}$ :
  - Inclusive  $dE/dx$  distribution fitted in 10 bands of  $\log[p_{TOT}/1\text{GeV}/c]$
  - Fit with 4 gaussian sum for  $\alpha = \pi, K, p, e$
  - Probability for a track with energy loss  $x_i$  of being a proton:
$$P = f^p(x_i, p_i) / (f^\pi(x_i, p_i) + f^K(x_i, p_i) + f^p(x_i, p_i) + f^e(x_i, p_i))$$





# Split tracks & the $q_{inv}$ cut

- Events may contain **split tracks**: sections of the same track erroneously identified as **a pair of tracks** that are close in momentum space.
- Intermittency analysis is based on pairs distribution  $\Rightarrow$  split tracks can create a **false positive**, and so must be **reduced** or **removed**.
- **Standard cuts** remove part of split tracks. In order to estimate the residual contamination, we check the  $q_{inv}$  distribution of track pairs:

$$q_{inv}(p_i, p_j) \equiv \frac{1}{2} \sqrt{-(p_i - p_j)^2},$$

$p_i$  : 4-momentum of  $i^{th}$  track.

- We calculate the ratio of  $q_{inv}^{data} / q_{inv}^{mixed}$ . A **peak** at low  $q_{inv}$  (below 20 MeV/c) indicates a possible split track contamination that must be removed.

# Comparison with simulated C+C – PHSD

- Simulation parameters:
  - ①  $^{12}\text{C}$  (beam) +  $^{12}\text{C}$  (target)
  - ② Beam energy: 150A GeV (target rest frame)
  - ③ Minimum bias  $\Rightarrow b_{\text{max}} \sim 6 \text{ fm}$
  - ④ Total number of simulated events: 100K
- Centrality selection by cut on minimum # of wounded nucleons:

$n_{\text{wounded}}^{\text{min}}$	#events	$\langle p \rangle_{ p_T  \leq 1.5 \text{ GeV},  y_{CM}  \leq 0.75}$ Non-empty	$\Delta p_{x,y}$ With empty	$\Delta p_{x,y}$
13	11684	$1.50 \pm 0.80$	$0.85 \pm 0.95$	0.48
14	10003	$1.54 \pm 0.80$	$0.90 \pm 1.00$	0.485
15	8245	$1.57 \pm 0.80$	$0.95 \pm 0.99$	0.48
16	6524	$1.60 \pm 0.80$	$0.99 \pm 1.00$	0.48
17	5008	$1.63 \pm 0.85$	$1.06 \pm 1.03$	0.48
18	3534	$1.70 \pm 0.90$	$1.11 \pm 1.11$	0.48

- NA49 “C” + C @ 158 GeV:  $\langle p \rangle = 1.6 \pm 0.9 \Rightarrow$  close to 10% most central C+C PHSD

# Simulating non-critical Be+Be – PHSD\*

- Simulation parameters:
  - ①  ${}^7\text{Be}$  (beam) +  ${}^9\text{Be}$  (target)
  - ② Beam energy: 150A GeV (target rest frame)
  - ③ Minimum bias  $\Rightarrow b_{\text{max}} \sim 6 \text{ fm}$
  - ④ Total number of simulated events: 74600
- Centrality selection by cut on minimum # of wounded nucleons:

$n_{\text{wounded}}^{\text{min}}$	#events	$\langle p \rangle_{ p_T  \leq 1.5 \text{ GeV},  y_{CM}  \leq 0.75}$ Non-empty	$\langle p \rangle_{ p_T  \leq 1.5 \text{ GeV},  y_{CM}  \leq 0.75}$ With empty	$\Delta p_{x,y}$
10	2632	$1.36 \pm 0.62$	$0.61 \pm 0.80$	0.48
11	1879	$1.39 \pm 0.64$	$0.65 \pm 0.82$	0.47
12	1219	$1.42 \pm 0.66$	$0.69 \pm 0.85$	0.48
13	710	$1.46 \pm 0.68$	$0.76 \pm 0.88$	0.47
14	358	$1.48 \pm 0.70$	$0.82 \pm 0.90$	0.44 - 0.48
15	108	$1.50 \pm 0.70$	$0.84 \pm 0.91$	0.43 - 0.49
16	21	$1.52 \pm 0.79$	$0.78 \pm 0.95$	0.40 - 0.53

PHSD data provided by Dr. Vitalii Ozvenchuk

\*[Cassing W. and Bratkovskaya E. L., Nucl. Phys. A **831** 215 (2009)]

# Simulating non-critical Be+Be – EPOS

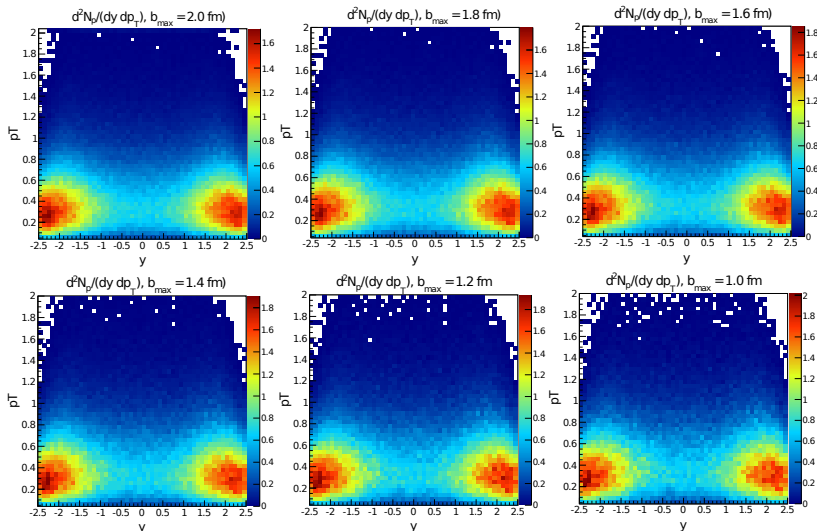
- Simulation parameters:

- ①  ${}^7\text{Be}$  (beam) +  ${}^9\text{Be}$  (target)
- ② Beam energy: 150A GeV (target rest frame)  $\Leftrightarrow \sqrt{s_{NN}} = 16.8$  GeV
- ③ Central collisions  $\Rightarrow b_{\text{max}} = 2.0$  fm
- ④ Total number of simulated events: 200K

$b_{\text{max}}$	#events	$\langle p \rangle_{ p_T  \leq 1.5 \text{ GeV},  y_{CM}  \leq 0.75}$ Non-empty	With empty	$\Delta p_{x,y}$
2.0	200,000	$1.41 \pm 0.69$	$0.66 \pm 0.85$	0.42
1.8	162,231	$1.43 \pm 0.70$	$0.69 \pm 0.87$	0.42
1.6	128,216	$1.44 \pm 0.71$	$0.72 \pm 0.88$	0.42
1.4	98,137	$1.46 \pm 0.73$	$0.74 \pm 0.90$	0.42-0.43
1.2	72,267	$1.47 \pm 0.73$	$0.76 \pm 0.91$	0.42-0.43
1.0	50,093	$1.48 \pm 0.74$	$0.78 \pm 0.92$	0.43

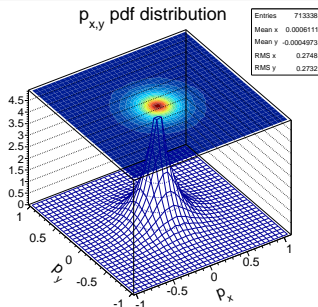
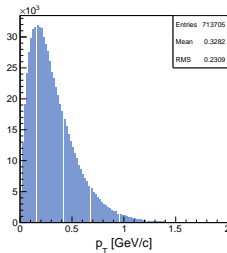
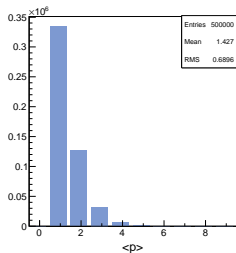
# Simulating non-critical Be+Be – EPOS, $y$ vs $p_T$

- $\langle p \rangle$  in mid-rapidity remains low, except for very central collisions



# Simulating critical Be+Be – CMC baryon

Parameter	$p_{\min}$ (MeV)	$p_{\max}$ (MeV)	$\lambda_{\text{Poisson}}$	$T_c$ (MeV)
Value	0.85	1200	0.76	163



# Simulating non-critical Ar+Sc – EPOS

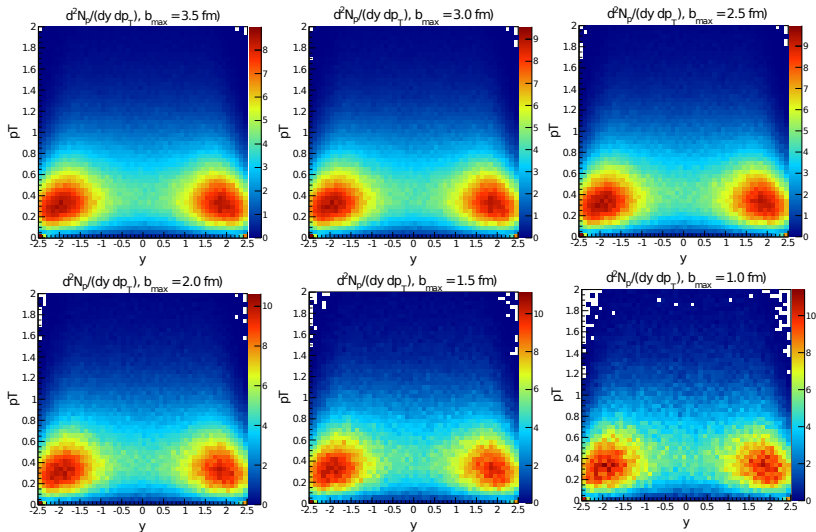
- Simulation parameters:

- ①  $^{40}\text{Ar}$  (beam) +  $^{45}\text{Sc}$  (target)
- ② Beam energy: 150A GeV (target rest frame)  $\Leftrightarrow \sqrt{s_{NN}} = 16.8$  GeV
- ③ Central collisions  $\Rightarrow b_{\text{max}} = 3.5$  fm
- ④ Total number of simulated events: 100K

$b_{\text{max}}$	#events	$\langle p \rangle_{ p_T  \leq 1.5 \text{ GeV},  y_{CM}  \leq 0.75}$		$\Delta p_{x,y}$
		Non-empty	With empty	
3.5	100,000	$5.3 \pm 2.5$	$5.3 \pm 2.4$	0.490
3.0	73,452	$5.6 \pm 2.5$	$5.7 \pm 2.5$	0.495
2.5	50,891	$5.9 \pm 2.5$	$6.0 \pm 2.5$	0.495
2.0	32,591	$6.2 \pm 2.5$	$6.2 \pm 2.5$	0.500
1.5	18,345	$6.4 \pm 2.6$	$6.5 \pm 2.6$	0.500
1.0	8,285	$6.6 \pm 2.6$	$6.5 \pm 2.6$	0.500
0.5	2,032	$6.7 \pm 2.7$	$6.8 \pm 2.7$	0.500

# Simulating non-critical Ar+Sc – EPOS, $y$ vs $p_T$

- $\langle p \rangle$  in mid-rapidity acceptable for  $b_{\max} \leq 3.5$





# Simulating critical Ar+Sc – CMC baryon

Parameter	$p_{\min}$ (MeV)	$p_{\max}$ (MeV)	$\lambda_{\text{Poisson}}$	$T_c$ (MeV)
Value	0.41	1200	5.3	163

